



# QFT I

## Exercise Sheet 3

**ETH**Eidgenössische Technische Hochschule Zürich  
Swiss Federal Institute of Technology Zurich

Lecturer: Prof. Thomas Gehrman  
Assistants: Pier Francesco Monni, Erich Weihs,  
Stephan Bühler, Mathias Ritzmann, Gabriel Abelof  
[www-theorie.physik.uzh.ch/~pfmonni/QFTI\\_HS10/](http://www-theorie.physik.uzh.ch/~pfmonni/QFTI_HS10/)

Issued : 04.10.2010  
Due : 11.10.2010  
Discussion : 15.10.2010

### Exercise 1 [*Chirality and helicity*]

Show that for  $m = 0$ , chirality and helicity are equivalent.

### Exercise 2 [*Angular momentum conservation*]

By explicit calculation, show that the total angular momentum

$$\vec{J} = \vec{x} \times \vec{p} + \frac{\hbar}{2} \vec{\Sigma}$$

and the Dirac-Hamilton-operator for a central potential

$$H = c \sum_{k=1}^3 \gamma^0 \gamma^k p^k + \gamma^0 m c^2 + e \Phi(|\vec{x}|), \quad \text{commute.}$$

### Exercise 3 [*Gordon identity*]

a) Derive the Gordon identity that holds for massless particles:

$$\langle u^\pm | \gamma^\mu | u^\pm \rangle := \bar{u}_\pm(k) \gamma^\mu u_\pm(k) = 2k^\mu .$$

Here, we have introduced the spinor helicity notation which is very useful to obtain compact expressions for multi-particle scattering amplitudes.

b) Prove that the projection operator onto states with positive/negative helicity can be expressed as

$$|u^\pm\rangle\langle u^\pm| = \frac{1}{2}(1 \pm \gamma_5) \not{k} .$$

*Hint:* First, prove that the massless spinors can be written in this useful form ( $c = 1$ ):

$$u_+(k) = v_-(k) = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{k^+} \\ \sqrt{k^-} e^{i\varphi_k} \\ \sqrt{k^+} \\ \sqrt{k^-} e^{i\varphi_k} \end{pmatrix}, \quad u_-(k) = v_+(k) = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{k^-} e^{-i\varphi_k} \\ -\sqrt{k^+} \\ -\sqrt{k^-} e^{-i\varphi_k} \\ \sqrt{k^+} \end{pmatrix}$$

with

$$e^{\pm i\varphi_k} = \frac{k^1 \pm ik^2}{\sqrt{(k^1)^2 + (k^2)^2}} = \frac{k^1 \pm ik^2}{\sqrt{k^+ k^-}}, \quad k^\pm = k^0 \pm k^3$$

**Exercise 4** [*Maxwell's equations*]

Starting from the Lagrangian of the electromagnetic field,

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - j_\mu A^\mu$$

derive the inhomogenous Maxwell's equations (the Euler-Lagrange equations of  $\mathcal{L}$ ). The field strength tensor and the four-current are given by

$$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu \equiv \begin{pmatrix} 0 & -E_1 & -E_2 & -E_3 \\ E_1 & 0 & -B_3 & B_2 \\ E_2 & B_3 & 0 & -B_1 \\ E_3 & -B_2 & B_1 & 0 \end{pmatrix}, \quad j^\mu = \begin{pmatrix} \rho \\ \vec{j} \end{pmatrix}.$$